

Home Search Collections Journals About Contact us My IOPscience

Positrons in gas filled traps and their transport in molecular gases

This article has been downloaded from IOPscience. Please scroll down to see the full text article. 2011 J. Phys.: Conf. Ser. 262 012046 (http://iopscience.iop.org/1742-6596/262/1/012046) View the table of contents for this issue, or go to the journal homepage for more

Download details: IP Address: 147.91.83.226 The article was downloaded on 06/07/2011 at 11:35

Please note that terms and conditions apply.

Positrons in gas filled traps and their transport in molecular gases

Z Lj Petrović¹, A Banković¹, S Marjanović¹, M Šuvakov¹, S Dujko^{1,2,3}, G Malović¹, R D White³ and S J Buckman⁴

¹Institute of Physics, University of Belgrade, Pregrevica 118, POB 68, Zemun, Serbia

²Centre for Mathematics and Informatics, POB 94079, 1090 GB Amsterdam, The Netherlands

³ARC Centre for Antimatter-Matter Studies, James Cook University, Townsville 4810, QLD, Australia

⁴ARC Centre for Antimatter-Matter Studies, The Australian National University, Canberra, ACT, 0200, Australia

zoran@ipb.ac.rs

Abstract. In this paper we give a review of two recent developments in positron transport, calculation of transport coefficients for a relatively complete set of collision cross sections for water vapour and for application of they Monte Carlo technique to model gas filled subexcitation positron traps such as Penning Malmberg Surko (Surko) trap. Calculated transport coefficients, very much like those for argon and other molecular gases show several new kinetic phenomena. The most important is the negative differential conductivity (NDC) for the bulk drift velocity when the flux drift velocity shows no sign of NDC. These results in water vapour are similar to the results in argon or hydrogen. The same technique that has been used for positron (and previously electron) transport may be applied to model development of particles in a Surko trap. We have provided calculation of the ensemble of positrons in the trap from an initial beam like distribution to the fully thermalised distribution. This model, however, does not include plasma effects (interaction between charged particles) and may be applied for lower positron densities.

1. Introduction

It has recently become possible to model positron transport in gaseous media based on the measured scattering cross sections. The efforts of groups in San Diego [1-2], Trento [3], London [4], Detroit [5] Swansea [6] and recently Canberra [7,8] led to measurements of a wide range of data which cover all relevant momentum and energy transfer processes as well as the processes that change the number of particles (annihilation and positronium formation). Sometimes, however, one needs to complete the set with theoretical results which are also becoming reasonably abundant [9-11], preferably those results based on theories that were tested against experiments on other similar targets. In any case sufficiently detailed and complete sets are now available for argon, hydrogen, nitrogen and few more relevant gases.



Figure 1. Rate coefficients for positrons in water vapour: Ps formation and inelastic losses.

Having in mind that the principal constituent of a human body is H₂O, we attempted to compile a set of cross sections that could be used to model behaviour of positrons in a human body. While a combination of experimental data and calculations was sufficient with some extrapolations to form the set, we have taken advantage of the recent measurements at the ANU in Canberra to provide two of the most important cross sections [8], the total scattering and the positronium formation cross section. It is important to note that all of the relevant processes, while providing a good balance and effective behaviour of the positron swarms, cannot be tested against the measured transport data yet as such experiments (with two exceptions [12,13])

are not available at the moment in a sufficiently wide range of normalized electric fields (E/N).

In our previous reviews of preliminary results we presented the first results for positron transport including the data for argon and some aspects of transport in molecular gases and in the crossed magnetic and electric fields [14]. Finally we have dealt with thermalization of positronium and transport in liquids which are associated with applications of positron transport [15]. Final results have been published for argon [16] and partly for nitrogen [18] with a special analysis of the negative differential conductivity (NDC) induced by positronium formation [16].

In this paper we present new developments with transport of positrons in water vapour and first attempts to model low density positron traps in gases.

2. Positron transport in water vapour

Before addressing the results we should perhaps discuss whether one has sufficient motivation for swarm (transport) studies of positrons and how those could be justified by applications.

The time scales of possible swarm experiments in gases that are at a high enough pressure are very short and yet positrons survive long enough to allow us determination of transport coefficients [12,13]. A review of past, existing and possible future swarm experiments with the discussion how to overcome formidable difficulties facing studies of positron swarms was given by Charlton [19]. It is important to note that a swarm system should not be viewed as a competition with traps in efficiency of storing swarms. Absence of multiplication (as compared to electrons) does not allow us to form stable ensembles of particles, so for any measurement the price is in the loss of particles. Such an experiment would thus provide transport data allowing us to adjust the desired properties by electric and magnetic fields- see accompanying paper [20], model some of the applications and also giving us a possibility to test and benchmark the calculations of the models of such applications. This is especially important as it would be very difficult to base all the comparisons on predicted trajectories and some averaged, accurately measured and calculated, data are needed. Finally such experiments would provide a natural occurrence of some of the predicted and yet to be predicted kinetic phenomena driven by non-conservative processes, Ps formation in this case.

Transport coefficients in water vapour calculated by a standard Monte Carlo procedure [16] show similar features as in argon and hydrogen. In other words, while mean energy is similar to what was observed for electrons, the drift velocity shows the duality of drift velocities, with NDC unlike anything seen for electrons (although the effect was predicted for electrons [21] in a model gas). The *flux* drift velocity increases gradually throughout the range of electric fields covered and at the same

Journal of Physics: Conference Series 262 (2011) 012046



kinetic energy of positrons in Surko trap.



Figure 3. Temporal development of the energy distribution function (probability of finding particles in a given energy bin) of positrons in a trap.

doi:10.1088/1742-6596/262/1/012046

time the *bulk* drift velocity shows a very pronounced NDC with a maximum difference between the two that is greater than an order of magnitude. With the exception of hydrogen there are no experiments confirming such drift velocities. Our good experience with electron swarms teaches us that our predictions are quite accurate but one would prefer to have the transport data to normalize the entire sets of cross sections.

One example of the data shown here in Fig. 1 is the rate for Ps formation compared to other elastic and inelastic processes.

3. Monte Carlo modelling of a Penning Malmberg Surko trap

The initial development and analysis of a Surko trap was done on the basis of assumed beam like properties of particles in the trap. Yet particles in such traps start from a beam like distribution and thermalize to an isotropic Maxwellian distribution at room temperature. In that respect the Surko trap is similar to the Frank Hertz experiment that will have its 100th anniversary in 2014. Very much like in the case of the Frank Hertz experiment [22] we believe that proper modeling of the Surko trap may be performed by swarm techniques that may provide description of a nonhydrodynamic particle kinetics (when the swarm is not in equilibrium with the local field but may depend on the spatial or temporal history).

In Figure 2 we show temporal development of the mean energy of positrons in trap. At the same time we show energy distributions at several times after the particles have been introduced to the trap. Unlike the real trap where particles are fed into the trap continuously here we introduce the whole batch together in order to follow the temporal development.

While providing the qualitative physical picture of the trap that is already available, the modeling also does so in a quantitative manner thus allowing comparisons between different gases, different fields, analysis of a possible benefit of added axial or radial electric fields.

Acknowledgements

This work was funded by the MNTRS project 141025, COE Nonequilibrium processes in plasmas and environmental science, ARC Centre for Antimatter-Matter Studies and the International Science Linkages Program. Authors are grateful to J. Marler, W. Tattersall, J. Sullivan and C. Makochekanwa for collaboration and advice and to G. Garcia, C. Surko and M. Charlton for many useful comments.

References

- [1] Marler JP and Surko CM 2005 *Phys. Rev. A* **72** 062702
- [2] Sullivan JP, Gilbert SJ, Marler JP, Greaves RG, Buckman SJ and Surko CM 2002 Phys. Rev. A 66 042708

Marler JP, Sullivan JP and Surko CM 2005 Phys. Rev. A 71 022701

- [3] Zecca A, Chiari L, Sarkar A, Lima MAP, Bettega MHF, Nixon KL and Brunger MJ 2008 Phys. Rev. A 78 042707
 Zecca A, Sanyal D, Chakrabarti M and Brunger MJ 2006 J. Phys. B: At. Mol. Opt. Phys. 39 1597
- [4] Laricchia G, Van Reeth P, Szluinska M and Moxom J 2002 J. Phys B. 35 2525
- [5] Kauppila WE, Stein TS, Smart JH, Dababneh MS, Ho YK, Downing JP and Pol V 1981 *Phys. Rev.* A 24 725
 Hoffman KR, Dababneh MS, Hsieh YF, Kauppila WE, Pol V, Smart JH, Stein TS 1982

Hoffman KR, Dababneh MS, Hsieh YF, Kauppila WE, Pol V, Smart JH, Stein TS 1982 Phys. Rev. A 25 1393

- [6] Charlton M 1985 *Rep. Prog. Phys.* 48 737
 Charlton M and Humberston JW 2001 *Positron Physics* (Cambridge University Press).
- [7] Sullivan J P, Makochekanwa C, Jones A, Caradonna P and Buckman S J, 2008 J. Phys. B: At. Mol. Opt. Phys. 41 081001
- [8] Makochekanwa C, Bankovic A, Tattersall W, Jones A, CaradonnaP, Slaughter D.S, Nixon K, Brunger M.J., Petrovic ZLj, Sullivan J.P. and Buckman S.J., 2009 New Journal of Physics 11 103036
- [9] McEachran R P 2006 Private communication
- [10] Gilmore S, Blackwood JE and Walters HRJ 2004 Nucl. Instr. Meth. B 221 129
- [11] F.A. Gianturco, T. Mukherjee, Phys. Rev. A 55 (1997) 1044.
- [12] Bose N, Paul D. A. L., and Tsai J. S., 1981 J. Phys. B 14, L227.
- [13] Charlton, M 1985 J. Phys. B 18, L667
- [14] Marler JP, Petrović ZLj, Banković A, Dujko S, Šuvakov M, Malović G and Buckman SJ 2009 *Phys. Plasmas*
- [15] Sullivan J P, Buckman S J, Jones A, Caradonna P, Makochekanwa C, Slaughter D, Petrović Z Lj, Banković A, Dujko S, Marler J P and White R D 2009 Journal of Physics: Conference Series 162 (2009) 012002
- [16] Šuvakov M, Petrović ZLj, Marler JP, Buckman SJ, Robson RE and Malović G 2008 New J. Phys. 10 No 5 053034
- [17] Banković A, Petrović ZLj, Robson RE, Marler JP, Dujko S and Malović G 2008 Nucl. Instrum.Methods B
- [18] Banković A, Marler JP, Šuvakov M, Malović G and Petrović ZLj 2008 *Nuclear Inst. Methods, B* 266 462
- [19] Charlton M, Journal of Physics: Conference Series 162 (2009) 012003
- [20] Banković A, Dujko S, White RD, Malović G, Buckman SJ and Petrović ZLj, this issue
- [21] Vrhovac SB and Petrović ZLj 1996 Phys. Rev. E 53 4012
- [22] Dujko S, White R D and Petrović Z Lj, 2008 J. Phys. D: Appl. Phys. 41 025206